## Two-dimensional projections of a hypercube

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We present a method to project a hypercube of arbitrary dimension on the plane, in such a way as to preserve, as well as possible, the distribution of distances between vertices. The method relies on a Monte Carlo optimization procedure that minimizes the squared difference between distances in the plane and in the hypercube, appropriately weighted. The plane projections provide a convenient way of visualization for dynamical processes taking place on the hypercube.

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Dynamical models where the state of a system is represented by an ordered array of binary variables are ubiquitous in statistical physics, especially in its interdisciplinary applications. Perhaps the widest class of models that admit such Boolean-like representation is constituted by binary cellular automata [1,2]. Specific applications include biological evolution at the levels of molecules [3–5], cells [6], individuals [7], and species [8–10], as well as social and socioeconomical behavior [11,12]. Moreover, genetic algorithms are typically applied to systems whose configuration is described by means of binary sequences [13]. These models may involve large populations of interacting agents, each of them described as a time-dependent array of bits, which requires assigning an evolving density to each possible binary sequence.

While the configuration space of a binary sequence of length L is naturally represented as the set of  $2^{L}$  vertices of an L-dimensional hypercube, its visualization can be disappointingly difficult, even for L not very large. On the other hand, besides a quantitative characterization of the system dynamics through its collective properties, it is sometimes desirable to rely on a geometrical depiction where the dynamics can be followed, for instance, on the computer screen. The purpose of this paper is to present a method to project the vertices of a hypercube of arbitrary dimension onto a set of points in the plane, with the condition of preserving (as much as possible) the structure of the distance distribution on the hypercube. The motivation of this condition is that many dynamical processes depend on the Hamming distance-i.e., the number of different bits-between binary sequences, and we require this feature to be well represented by the Euclidean distance between the corresponding points in the plane projection.

Let  $h_{ij}$  be the Hamming distance between vertices *i* and *j* in the hypercube, and  $d_{ij}$  the Euclidean distance between points *i* and *j* in their plane projection. We define the function

$$E = \sum_{i,j} (d_{ij} - h_{ij})^2, \qquad (1)$$

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which characterizes how different are the distances between pairs of vertices and their projections. Our goal is to find a plane distribution that minimizes E, thus optimizing the plane representation of the hypercube with respect to the distance between pairs. We have implemented a Monte Carlo method to stochastically approach the optimal solution—the configuration of minimum "energy" E. Starting from a random initial configuration on the plane, each point performs a walk with fixed step length r and directions chosen at random with uniform probability in  $[0,2\pi)$ . Each step of this walk produces a change in the configuration and, hence, in the distances  $d_{ij}$ , which implies a variation  $\Delta E$  in the energy. The new configuration is accepted with probability

$$p = \begin{cases} \exp(-\Delta E/T) & \text{if } \Delta E > 0, \\ 1 & \text{otherwise,} \end{cases}$$
(2)

and rejected with probability 1-p. The "temperature" T parametrizes this probability and allows the usual implementation of a simulated annealing, where the procedure starts with a high temperature that enables the system to explore a wide range in configuration space. Progressively, the temperature is reduced and the system freezes in one of the many local minima of the energy, typically not far away from the global minimum if the annealing is made slowly enough.

We have carried out the described procedure both interactively, reducing by hand the temperature while monitoring the configuration of the system on the computer screen, and automatically, by implementing a programmed reduction of the temperature. Our experiments show that essentially the same state is achieved in almost all the realizations. This implies that the energy landscape, while rugged, does not posses deep local minima that could capture the configuration far from the optimal one. The typical final configuration for L = 10 (N = 1024 points) is shown in Fig. 1(a). The selfsimilarity of its structure is remarkable, since no such property is present in the hypercube. Despite the appeal that this self-similar projection may have, it turns out that such projection is not well suited for our purpose. Vertices that are relatively near in the hypercube result rather far away in the projection. As an illustration, the first neighbors  $j(h_{ii}=1)$  of a given vertex *i* are shown in the figure. It is apparent that the Euclidean distance of some of them from the reference ver-



FIG. 1. Two plane projections of a ten-dimensional hypercube: (a) using the energy defined in Eq. (1); (b) using the energy defined in Eq. (3). Lines join the projections of a randomly chosen vertex of the hypercube and its nearest neighbors.

tex is comparable with the size of the system. Moreover, many other vertices which should be farther away from vertex i are, in fact, much closer.

One way to solve this difficulty is to modify the definition of E, such that near neighbors have more weight than distant neighbors. In fact, the energy (1) overemphasizes the effect of large distances. The weights can be chosen in a variety of ways, even as new tunable parameters in the optimization process [5]. We have implemented the following simple alternative, with fixed weights,

$$E = \sum_{i \neq j} \left( \frac{d_{ij} - h_{ij}}{h_{ij}} \right)^2.$$
(3)

The final configuration, which we will term "homogeneous," is shown in Fig. 1(b). Neighbor vertices of a vertex i now result mapped onto points that surround the point i, which makes this projection much more satisfying. Certainly, however, some vertices result mapped near the border of the



FIG. 2. Distribution of distances in the hypercube, the plane projections, and selected subsets of these: (a) self-similar; (b) homogeneous.

circle, and the arrangement of their neighbors is slightly different than that of vertices mapped in the middle of the set. We analyze below how this affects the distribution of distances.

A good characterization of the projections, and a quantitative way for comparing them, is the distribution of distances in each set. In the hypercube the distribution of distances to a vertex is the same for every vertex, and in fact is analytically found to be a binomial distribution. In the twodimensional projections there is a different distribution for each point of the set. In Figs. 2(a) and 2(b), we show normalized distributions of distances for the self-similar and the homogeneous projections, respectively. In both figures, the black circles show the distribution of the distance to any vertex in the hypercube. Even though the distances form a discrete set, we show lines connecting the points to ease the reading of the graph. The other three curves shown in each plot correspond to the plane projections. The black squares correspond to an average over all the points in the sets. Triangles show averages performed on either the 10% of the points that form the external corona of the projection, or the 10% of its more central points. For the points of these subsets, still, all the distances to other points of the whole set are taken into account in the distributions.

The most immediate observation regarding Fig. 2 is the difference between the distributions in the two projections. The self-similar projection displays rugged distributions that reflect the hierarchical geometrical arrangement of the points. In the homogeneous projection, instead, the distributions are smooth, as in the hypercube. To this extent, the homogeneous projection can be said to represent more accurately the distribution of distances present in the hypercube.

The distribution averaged over the whole set appears, however, skewed towards smaller distances, with a maximum around d=3, instead of the most represented distance h=5 of the hypercube. Interestingly, Fig. 2(b) shows that the *outer* 10% points considerably correct this skew. In other words, a point near the border of the circular array of the projection has a distribution of distances to the other points in the set, which is rather similar to the distribution of a vertex of the hypercube.

An appraisal of the plane projections of the hypercube in a dynamical context results from their application of a diffusion process. Let us suppose that, at each time step, a random walker jumps from a vertex of the hypercube to one its neighbors with equal probability. The average distance Dfrom the initial site, as a function of time, is shown in Figs. 3(a) and 3(b) as black circles. The inset in both figures displays the same curve in double logarithmic scales, showing an initial behavior of the form  $D(t) \sim t^{1/2}$ , like in a regular random walk in Euclidean space, followed by a saturation as the hypercube space is fully explored. The average distance as measured on the plane is shown in Figs. 3(a) and 3(b) for the self-similar and the homogeneous projections, respectively. As expected from the distance distribution discussed above, the results for the plane projections depend on whether the initial point of the walker is at the border or at the center of the set. These two cases are shown in Figs. 3(a)and 3(b) as triangles pointing upward and downward, respectively. From this dynamical point of view, interior points behave equally bad in both projections. The most faithful representation of the process in a plane projection is the one given by one of the border points of the homogeneous set [Fig. 3(b), up triangles]. Diffusion starting at these points behaves similarly as from points of the hypercube, both in the short and in the long time regimes, as seen in the linear and the logarithmic plots.

Our main goal of obtaining a sensible plane projection of the hypercube with the purpose of visualizing a dynamical process has been achieved, to an acceptable extent, by the homogeneous projection. Suppose that a dynamical phenomenon is taking place in a neighborhood of vertex P of a hypercubical phase space. We need to build a homogeneous



FIG. 3. Average displacement as a function of time for diffusion in the hypercube and its plane projections. For the plane projections, diffusion starting from a point in the border and a point near the center are shown separately: (a) self-similar; (b) homogeneous. The straight lines in the insets have a slope of 1/2.

projection that maps vertex P to a point at the border of the plane set. This is easily done by generating a projection at random and identifying one of the points at the border first. Suppose that one such point is Q. Then, each vertex I of the hypercube is mapped to a point in the plane projection as

$$I \to (I \oplus P) \oplus Q, \tag{4}$$

where  $\oplus$  stands for the bitwise exclusive-OR (XOR) operator. The projection obtained in this way provides a nice plane visualization substrate for the process.

- [1] S. Wolfram, Rev. Mod. Phys. 55, 601 (1983).
- [2] S.A. Kauffman, *The Origins of Order* (Oxford University Press, Cambridge, 1993).
- [3] M.V. Volkenstein, *Physical Approaches to Biological Evolution* (Springer-Verlag, Berlin, 1994), Chap. 8.
- [4] M.A. Jiménez Montaño, C.R. de la Mora-Basáñez, and T. Pöschel, BioSystems 39, 117 (1996).
- [5] A. Coghlan, D.A. Mac Dónaill, and N.H. Buttimore, Bioinformatics 17, 676 (2001).
- [6] S.A. Kauffman, J. Theor. Biol. 22, 437 (1969); Nature (London) 244, 177 (1969).

- [7] T.J.P. Penna, J. Stat. Phys. 78, 1629 (1995).
- [8] S.A. Kauffmann and S. Johnsen, J. Theor. Biol. **149**, 467 (1991).
- [9] K. Kaneko and T. Ikegami, Physica D 56, 406 (1992).
- [10] R.V. Solé and S.C. Manrubia, Phys. Rev. E 54, R42 (1996).
- [11] R.G. Palmer, W.B. Arthur, J.H. Holland, B. LeBaron, and P. Tayler, Physica D 78, 1629 (1995).
- [12] G. Weisbuch, G. Deffuant, F. Amblard, and J.-P. Nadal, Complexity 7, 55 (2002).
- [13] J.H. Holland, Sci. Am. 267(1), 44 (1992).